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ACOUSTIC THERMOMETRIC MEASUREMENTS OF PROPELLANT GAS TEMPERATURES IN GUNS

Edward M. Schmidt
Edmund J. Gion
Donald D. Shear

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TABLE OF CONTENTS

	<u>Page</u>
LIST OF ILLUSTRATIONS	5
I. INTRODUCTION	7
II. INSTRUMENTATION AND EXPERIMENTAL DATA	9
III. DATA REDUCTION AND PRESENTATION	11
IV. CONCLUSIONS	14
REFERENCES	24
LIST OF SYMBOLS	25
DISTRIBUTION LIST	27

LIST OF ILLUSTRATIONS

<u>Figure</u>		<u>Page</u>
1.	Shadowgraph of Muzzle Blast	16
2.	Idealized, One-Dimensional Exit Flow	17
3.	Ratio of Energy Efflux Rates Versus Incident Flow Mach Number	17
4.	In-bore Wave Diagram	18
5.	Schematic of Test Set-up	18
6.	Muzzle Adaptor	19
7.	Sample Pressure Traces at $u_m = 366$ m/s	20
8.	Pressure Traces from Gauge 3 at Various Launch Velocities	21
9.	Wave Diagram of Muzzle Flow	22
10.	Measured Speed of Sound	22
11.	Measured Exit Pressure	23
12.	Measured Exit Temperature	23

I. INTRODUCTION

The muzzle blast generated during weapon firing, Figure 1, is driven by the expanding propellant gases ejected subsequent to projectile launch. It is of interest to model the blast field both to determine its environmental signature and its effect upon the projectile trajectory. Any quantitative model of this flow field requires the specification of accurate initial conditions, i.e., propellant gas properties at the muzzle during gun tube emptying. The propellant gas properties predicted by interior ballistics models¹⁻³ are dependent upon the method of treating the combustion and gas dynamic processes occurring within the gun tube. In general, the validity of these models is gauged against the accuracy with which they predict projectile acceleration and in-bore pressure histories. These parameters are not sufficient for muzzle blast calculations.

A propellant gas property of particular importance is temperature. The development of the muzzle blast has been shown^{4,5} to obey the predictions of spherical, strong blast theory⁶ based upon the initial rate of energy efflux from the muzzle. Gas temperature, through its impact upon the speed of sound, has a significant influence upon the initial rate of energy efflux. This may be demonstrated by considering the one-

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1. *Engineering Design Handbook, Ballistics Series, Interior Ballistics of Guns*, AMCP 706-150, U. S. Army Materiel Command, February 1965.
 2. P. G. Baer and J. M. Frankle, "The Simulation of Interior Ballistic Performance of Guns by Digital Computer Program," BRL Report No. 1183, U. S. Army Ballistic Research Laboratories, Aberdeen Proving Ground, Maryland, December 1962. AD 299980.
 3. A. Celmins, "Theoretical Basis of the Recoilless Rifle Interior Ballistic Code, RECRIF," BRLR (to be published), Ballistic Research Laboratories, Aberdeen Proving Ground, Maryland.
 4. J. I. Erdos and P. Del Guidice, "Calculation of Muzzle Blast Flow-fields," AIAA J., Vol. 13, No. 8, August 1975, pp. 1048-1055.
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dimensional flow of an ideal gas. The energy flux past any station at a given time is

$$\begin{aligned}\dot{E} &= \rho u \left(c_v T + \frac{u^2}{2} \right) A \\ &= \frac{\rho u}{\gamma - 1} \left(1 + \frac{\gamma(\gamma - 1)}{2} \left(\frac{u}{a} \right)^2 \right) A.\end{aligned}\quad (1)$$

Neglecting the projectile presence and two dimensional effects, the idealized flow geometry shown in Figure 2 is used to approximate the arrival of the propellant gases at the muzzle and subsequent expansion of these gases once the muzzle is breached. Using one-dimensional, unsteady flow theory⁷, an expression for the ratio of incident rate of energy flux, \dot{E}_1 , to rate of energy passing the muzzle, \dot{E}^* , may be derived which is only a function of the ratio of specific heats and incident Mach number, u_1/a_1 . For $\gamma = 1.25$, this ratio is plotted versus u_1/a_1 in Figure 3. The energy flux ratio, \dot{E}^*/\dot{E}_1 , varies strongly with incident Mach number in the range $0 < u_1/a_1 \leq 0.6$. As the sonic condition is approached, the energy flux ratio asymptotes toward unity due to decreasing strength of the muzzle expansion. For sonic and supersonic values of u_1/a_1 , there is obviously no return expansion from the muzzle. Since several important classes of weapons (e.g.: mortars, howitzers) launch projectiles in the low incident Mach number regime, it is of interest to determine the detailed exit properties under these conditions.

A variety of techniques has been applied to the measurement of propellant gas temperature during the interior ballistic cycle. Thermocouples⁸ have been mounted to the inner wall of the gun tube in attempting to directly measure gas temperature. Since flush mounted devices are immersed in the propellant gas boundary/thermal layer, thermocouple output does not reflect temperature in the inviscid main stream. Optical techniques such as two-color pyrometry⁹ and

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7. A. H. Shapiro, The Dynamics and Thermodynamics of Compressible Fluid Flow, Vol. II, Ronald Press, N.Y., 1954.
 8. G. S. Fulcher, "The Temperature of the Bore Surface of Guns," NDRC R A-201, National Defense Research Committee, July 1943.
 9. F. C. Kracek and W. S. Benedict, "An Experimental Study of Powder Gas Radiation and Temperature," NDRC R A-252, National Defense Research Committee, February 1944.

spectroscopy^{10,11} are also complicated by the presence of the wall boundary layer. Additionally; contamination of optical surfaces by particulate deposition during the measurement period further degrade these approaches. Klingenberg and Mach¹¹ present a calibration and data acquisition procedure which reduces the influence of this deposition on spectroscopic measurements; however, the effect of the wall thermal layer remains to be assessed.

In the present report, a novel method to obtain in-bore temperature measurements is presented. The approach is similar to that used in acoustic thermometry¹²; namely, the velocity of propagation of an acoustic wave is determined by measuring the time required for it to travel a known distance. Unlike acoustic thermometry in which an external signal is introduced into the test medium, the present technique measures the propagation velocity of an isentropic wave which is generated in the flow itself, the returning muzzle expansion, Figure 4. Pressure transducers, flush mounted in the tube wall, are used to sense passage of the muzzle expansion. Since pressure signals are impressed across the wall boundary layer, its affects upon the measurements are minimized. The details of the instrumentation and data reduction procedures applied to a 20mm gun will be presented in the subsequent sections.

II. INSTRUMENTATION AND EXPERIMENTAL DATA

A schematic of the test set-up is shown in Figure 5. Data is taken on the launch of a standard, 98 gram, training projectile from a 20mm gun. The gun has a length of 152cm, chamber volume of 41.7cm³, and twist of one turn in 25 calibers. A summary of launch velocities obtained by loading the cartridge with different weights of WC870 powder is shown in Table I. Launch velocity was measured from flash X-rays taken at three stations covering the first 60cm of projectile flight.

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10. G. Klingenberg and H. Mach, "Spektroskopische Temperaturbestimmung im unmittelbaren Mündungsbereich Eines Gewehres," ABF R 4/74, Arbeitsgruppe für Ballistische Forschung, Weil am Rhein, Germany, December 1974.
 11. G. Klingenberg and H. Mach, "Experimental Study of Non-Steady Phenomena Associated with the Combustion of Solid Gun Propellants," presented at the Sixteenth International Symposium on Combustion, Boston, Massachusetts, August 1976.
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$M_c (\text{Kg} \times 10^3)$	3.56	5.76	8.10	10.04	12.63	18.15	22.02	25.93	33.33
$u_m (\text{m/s})$	265	325	385	450	540	605	680	765	910

TABLE I: LAUNCH CONDITIONS

A muzzle adaptor, having an overall length of 8.6cm, Figure 6, was fabricated to accommodate and shock mount the array of pressure transducers. The adaptor bore diameter is 21mm to permit projectile clearance of the gauge surfaces. Since the weapon groove diameter is 20.75mm, there is a 0.25mm clearance between the projectile rotating band and the adaptor bore surface. Including the contribution of the nine lands, there is a 4.3 percent increase in cross sectional area between the gun and adaptor bores. Assuming a sonic flow is generated in the annulus between the projectile and adaptor surfaces, the resulting change in the flow properties behind the projectile may be approximated using steady flow theory. This calculation shows the temperature and pressure decrease by 0.01 percent and 0.04 percent, respectively. Such property variations are not significant in the current experiments which are intended to prove the viability of the measurement technique. The active transducers are three Kistler, 603A, piezo-electric, pressure transducers spaced at 2.5cm intervals from the front face of the device. The outputs of the transducers are recorded on Tektronix, Type 551, Dual-Beam Oscilloscopes. A uniform time base between traces is provided by triggering all scopes simultaneously from the output of gauge 1 as the projectile passes. Additionally, the discharge of the first X-ray tube is placed on each oscilloscope record by chopping the lower beam. In this manner, the projectile location may be related to the measured pressure variations.

A sample set of pressure traces obtained for a launch velocity of 366 m/s is shown in Figure 7. As the projectile passes gauge 1, all scopes are triggered simultaneously. When the projectile passes a gauge location, there is a rapid pressure rise as the gauge becomes exposed to the high pressure propellant gases. Naturally, gauge station 1 is exposed first followed by stations 2 and 3. As the projectile moves down the tube, in-bore waves reflected from the breech and the projectile base pass the gauge location. This results in a gradual decay of the measured pressure at the gauge station.

When the projectile clears the muzzle, the propellant gases expand into the surrounding atmosphere. Simultaneously, an expansion propagates into the gun tube bringing the muzzle to a choked or sonic condition. As the expansion passes the gauge station, the pressure begins to decrease more rapidly. This is clearly observed as the sudden change in slope on the pressure traces. The wave is propagating into the gun tube; thus, it arrives first at station 3 and then moves

to stations 2 and 1. The difference in arrival times is evident from the three traces shown in Figure 7. Since station 3 is closest to the center of the muzzle expansion, the rate of change of properties occurs very rapidly. As the muzzle expansion propagates further into the gun, it spreads out; thus the change in slope of the pressure traces is less severe. It is desirable to locate the pressure gauges as near to the muzzle as possible in order to minimize property variation due to in-bore waves and to provide a strong signal upon passage of the muzzle expansion. This is particularly true at high launch velocities, Figure 8. As the launch velocity increases, pressure decay due to in-bore waves arriving prior to the muzzle expansion becomes more severe. Simultaneously, the strength of the muzzle expansion decreases. In the limiting case, the launch velocity increases to a point where the propellant gas reaches a sonic condition prior to shot ejection. When this occurs, there is no muzzle expansion into the gun tube.

III. DATA REDUCTION AND PRESENTATION

A wave diagram of the near muzzle region is shown in Figure 9. The goal of this technique is to determine the velocity of propagation of the lead wave of the muzzle expansion, m-f-d-b. The instantaneous velocity of a left running wave is

$$c = \frac{dx}{dt} = u - a, \quad (2)$$

where u and a are the local flow velocity and speed of sound, respectively. The most direct means to obtain an estimate of the speed of sound is to assume the flow velocity is constant and equal to the launch velocity, u_m , evaluate the wave speed using the measured arrival times at the gauge stations, and invert Equation (2);

$$a = u_m - \frac{x_3 - x_1}{t_f - t_b}. \quad (3)$$

This approach does not account for property variations occurring prior to arrival of the muzzle expansion as evidenced by pressure decay in Figure 8.

A correction procedure may be developed from one-dimensional flow theory⁷ if the flow is assumed to be locally isentropic within test region, a-m-b. The following relation describes the property variations along a left running characteristic or within a region of right running simple waves:

$$\frac{a}{a_1} = 1 + \frac{\gamma-1}{2} \left(\frac{u}{a_1} - \frac{u_1}{a_1} \right), \quad (4)$$

where the subscript, 1, refers to a suitable reference point. For isentropic flow,

$$\frac{p}{p_1} = \left(\frac{a}{a_1} \right)^{\frac{2\gamma}{\gamma-1}} \quad (5)$$

These relations apply along the lead expansion wave, m-f-d-b; however, there is no pressure gauge exactly at the muzzle. This restricts the location of the initial reference point to be point e. The pressure-time history along the path e-f-d-b may be obtained from the measured data (pressures from e-f come directly from the trace of the third gauge, pressures at d and b come from the intercept values on traces 2 and 1).

If the projectile velocity is assumed to remain constant over the last 2.5cm of travel through the device, $u_e = u_m$, and if the path e-f is approximated as a simple wave region, Equation (4) may be used to derive the following relation:

$$\frac{u}{a_e} - \frac{a}{a_e} = \frac{u_m}{a_e} + \frac{3-\gamma}{\gamma-1} \frac{a}{a_e} - \frac{2}{\gamma-1} \quad (6)$$

Substituting Equation (5),

$$\frac{u}{a_e} - \frac{a}{a_e} = \frac{u_m}{a_e} + \frac{3-\gamma}{\gamma-1} \left(\frac{p}{p_e} \right)^{\frac{\gamma-1}{2\gamma}} - \frac{2}{\gamma-1} \quad (7)$$

From Equation (2), the incremental displacement, dx, of the lead wave in time, dt, is

$$\begin{aligned} dx &= (u-a) dt \\ &= \left(u_m + \frac{3-\gamma}{\gamma-1} \left[\frac{p}{p_e} \right]^{\frac{\gamma-1}{2\gamma}} a_e - \frac{2a_e}{\gamma-1} \right) dt, \end{aligned} \quad (8)$$

This expression may be integrated over the path of the wave, f-d-b, and inverted to solve for a_e .

$$a_e = (\Delta x - u_m \Delta t) / \left(\frac{3-\gamma}{\gamma-1} \int_{t_f}^{t_b} \left[\frac{p}{p_e} \right]^{\frac{\gamma-1}{2\gamma}} dt - \frac{2}{\gamma-1} \Delta t \right) , \quad (9)$$

where $\Delta x = x_b - x_f$,

$\Delta t = t_b - t_f$.

All of the terms in Equation (9) may be obtained from the pressure traces. The resulting correction to the speed of sound is small for low launch velocities; however, at 800 m/s, the value of a_e obtained from Equation (9) is fifteen percent lower than the speed of sound obtained from direct differencing, Equation (3).

The variation in speed of sound with launch velocity obtained by this method is presented in Figure 10. For launch velocities between 200 and 600 m/s, the speed of sound of the propellant gas prior to shot ejection remains constant at 675 m/s with a scatter of plus or minus 20 m/s in the data. Subsequent to the 600 m/s launch velocity, the propellant gas speed of sound begins to increase rapidly reaching a sonic value at a launch velocity of 780 m/s. Further variations in the speed of sound with launch velocity may not be measured using the present technique.

The pressure measured at gauge station 3 immediately after projectile passage is plotted as a function of launch velocity in Figure 11. The muzzle pressure shows an almost linear variation with velocity until 600 m/s when it increases suddenly. After the discontinuous jump, the pressure again increases linearly. This sudden increase in muzzle pressure occurs at the same launch velocity at which the speed of sound begins to increase; thus, suggesting the possibility of a change in the combustion process. The design launch velocity of the round is 1000 m/s. Lower velocities were obtained by simply removing propellant from the cartridge case without reducing the chamber volume. The propellant was held against the ignitor in the rear of the cartridge by a thin diaphragm. The free volume thus generated between the propellant bed and the base of the projectile may permit the establishment of pressure waves affecting the subsequent combustion history of the propellant.

The propellant gas temperature is computed from the measured sound speed by assuming the applicability of an equation of state. In the present case, it is assumed that by the time the propellant gases expand to the muzzle, they may be treated as an ideal gas. The

muzzle temperature may be evaluated from

$$T = \frac{a^2}{\gamma R} , \quad (10)$$

where¹

$$\gamma = 1.25,$$

$$R = 365.5 \text{ m}^2/\text{s}^2 \text{ } ^\circ\text{K} .$$

The temperature evaluated from the measured sound speed is shown in Figure 12. At low launch velocities, the muzzle temperature remains roughly constant with a value of 1000°K. Subsequent to 600 m/s, the temperature increases to a maximum value of 1320°K at 800 m/s.

The present measurements are compared with the predictions of a closed form solution¹ and numerical integration² of the interior ballistics of the 20mm gun. The closed form solution, based on the Mayer-Hart System, requires input of empirical values of projectile velocity and muzzle pressure for each launch condition. The numerical solution requires adjusting of the friction and burning rate coefficients to produce a match to the measured muzzle pressure and launch velocity at 1000 m/s. Subsequent solutions are extrapolated from this match condition. In both cases, the theoretical solutions predict temperature variations with launch velocity which are opposite to the measured trends. Agreement between theory and experiment occur only at a single intercept point. It should be noted that the "wiggle" in the Mayer-Hart solution occurs due to its reliance upon empirical input which possess discontinuities, Figure 11. Also shown in Figure 12 is the result of a spectroscopic measurement by Klingenberg and Mach¹¹. Their measurement was obtain immediately behind a 20mm projectile as it passed an optical window located 120cm from the breech. This data point compares well with the results of the current technique taken at a station 158cm from the breech of a 20mm gun.

IV. CONCLUSIONS

A procedure is developed which permits the determination of propellant gas temperature immediately behind the projectile at shot ejection. Pressure gauges are used to observe the time of passage of the muzzle expansion at a known location. A differencing relation is presented to calculate the speed of sound of this wave. Incorporated in this relation are corrections to account for changes in speed of sound and flow velocity over the measurement period.

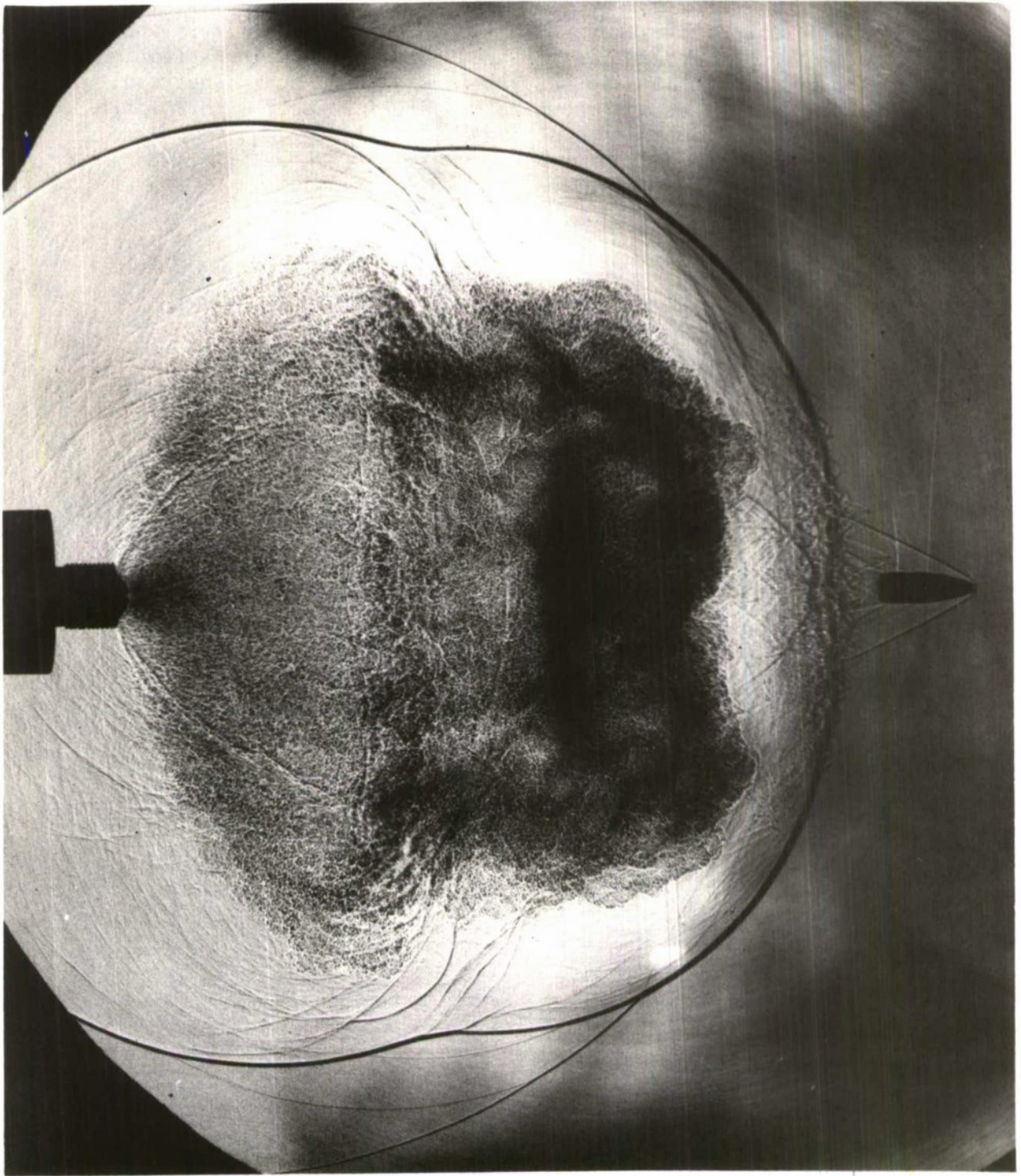


Figure 1. Shadowgraph of Muzzle Blast

The procedure is used to determine the propellant gas temperature at the muzzle of a 20mm gun over a range of launch velocities. Comparison of the measured temperature with simple interior ballistic analyses shows poor agreement indicating the need for continued investigation. Studies are being initiated to resolve this discrepancy. The results of the present technique are shown to compare favorably with spectroscopic temperature measurements obtained on a similar gun.

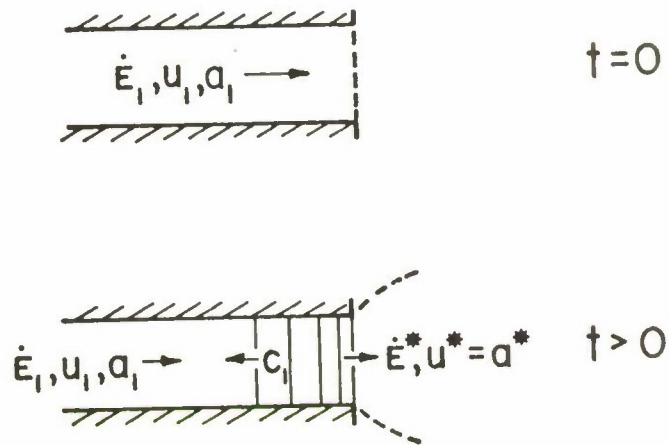


Figure 2. Idealized, One-Dimensional Exit Flow

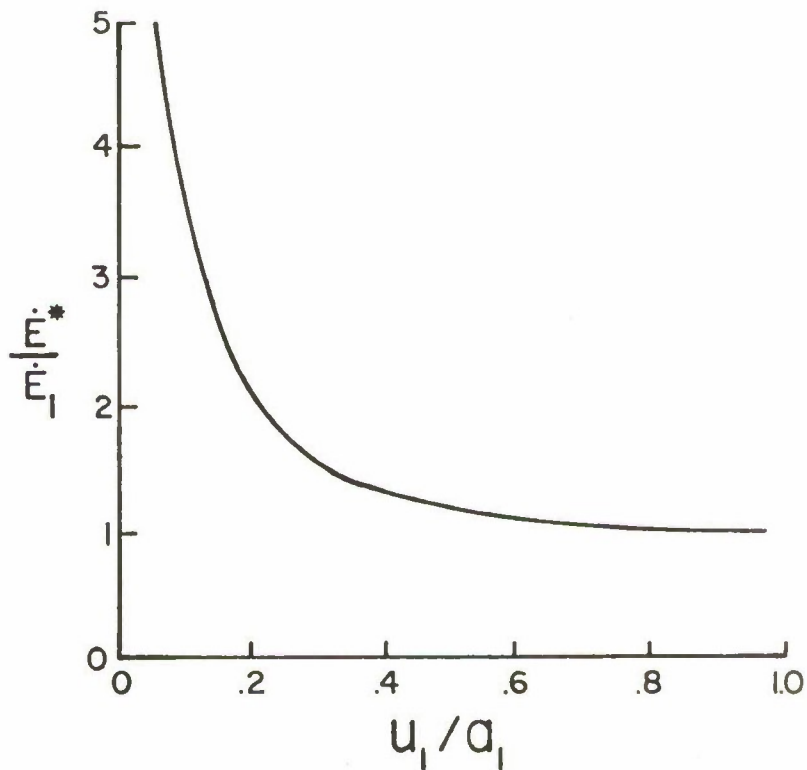


Figure 3. Ratio of Energy Efflux Rates Versus Incident Flow Mach Number

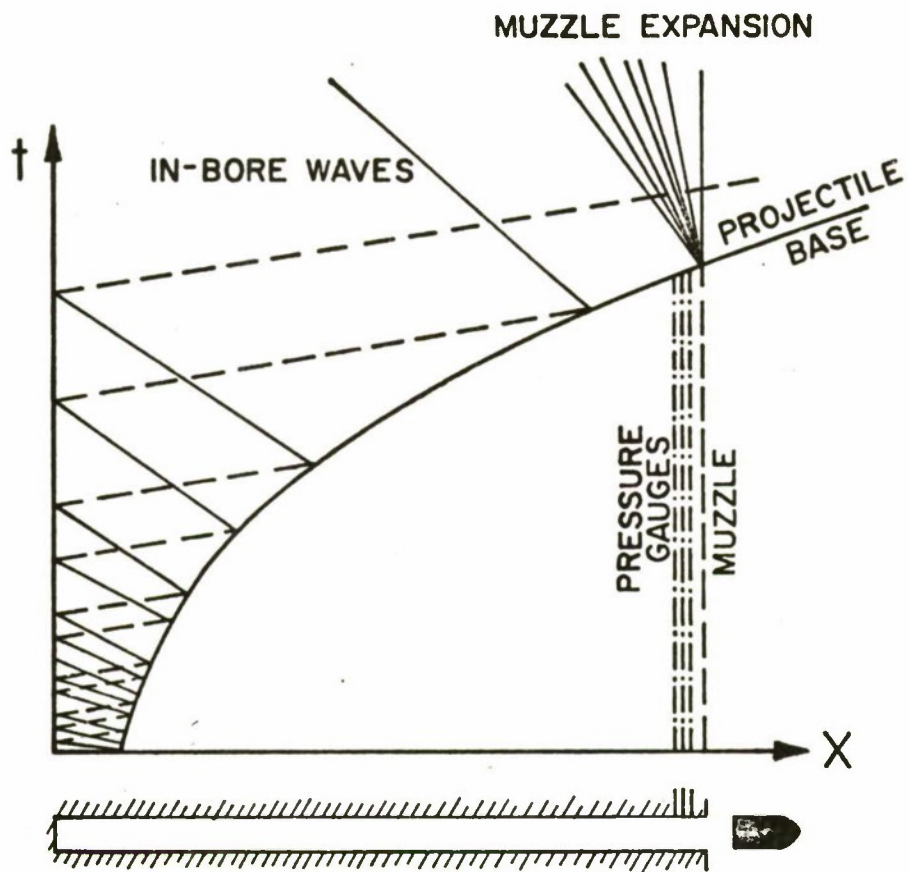


Figure 4. In-bore Wave Diagram

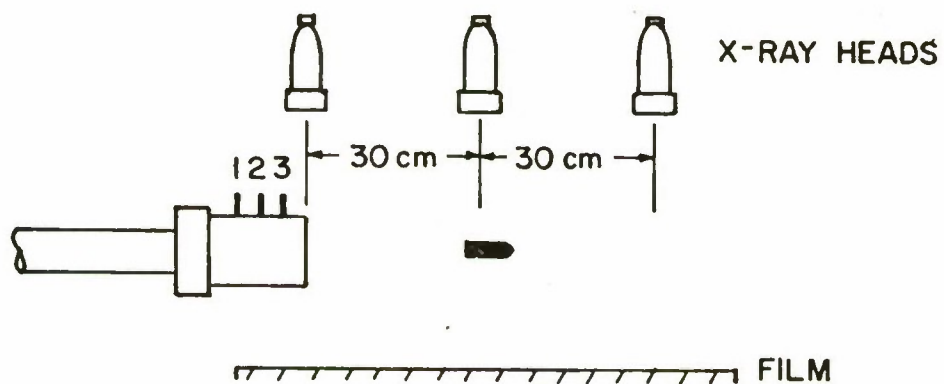


Figure 5. Schematic of Test Set-up

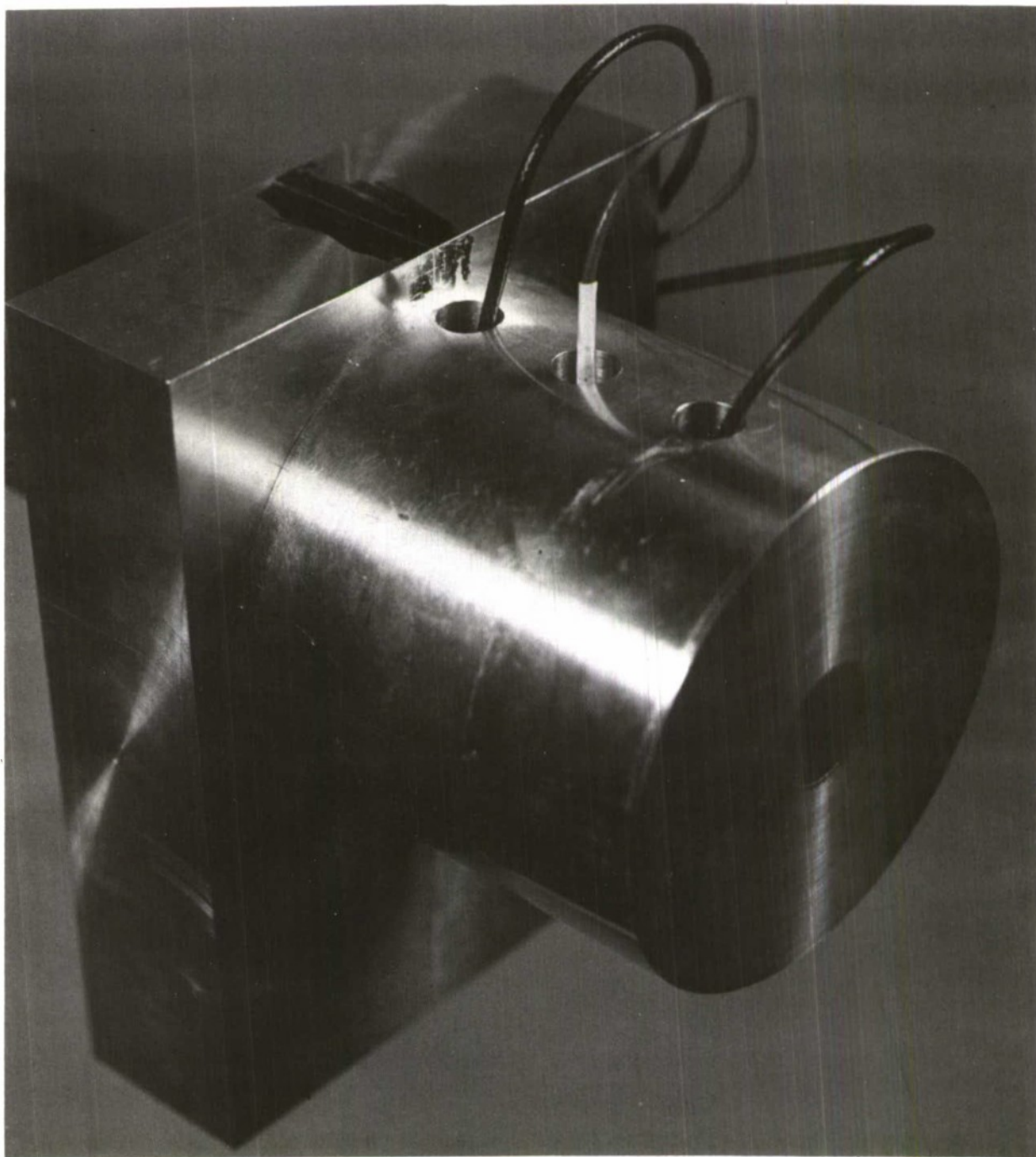
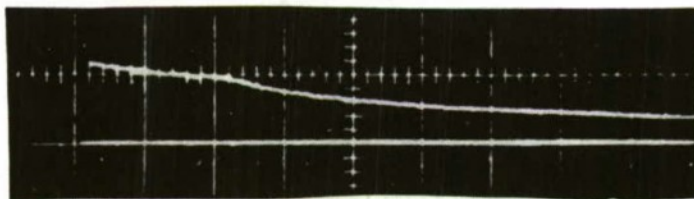
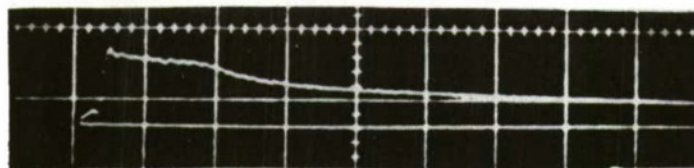


Figure 6. Muzzle Adaptor

STATION 1



STATION 2



STATION 3

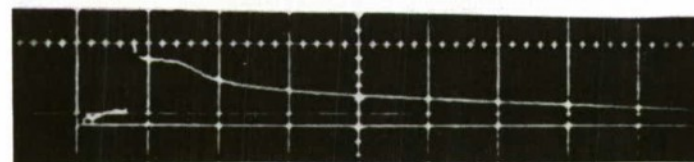


Figure 7. Sample Pressure Traces at $u_m = 366$ m/s

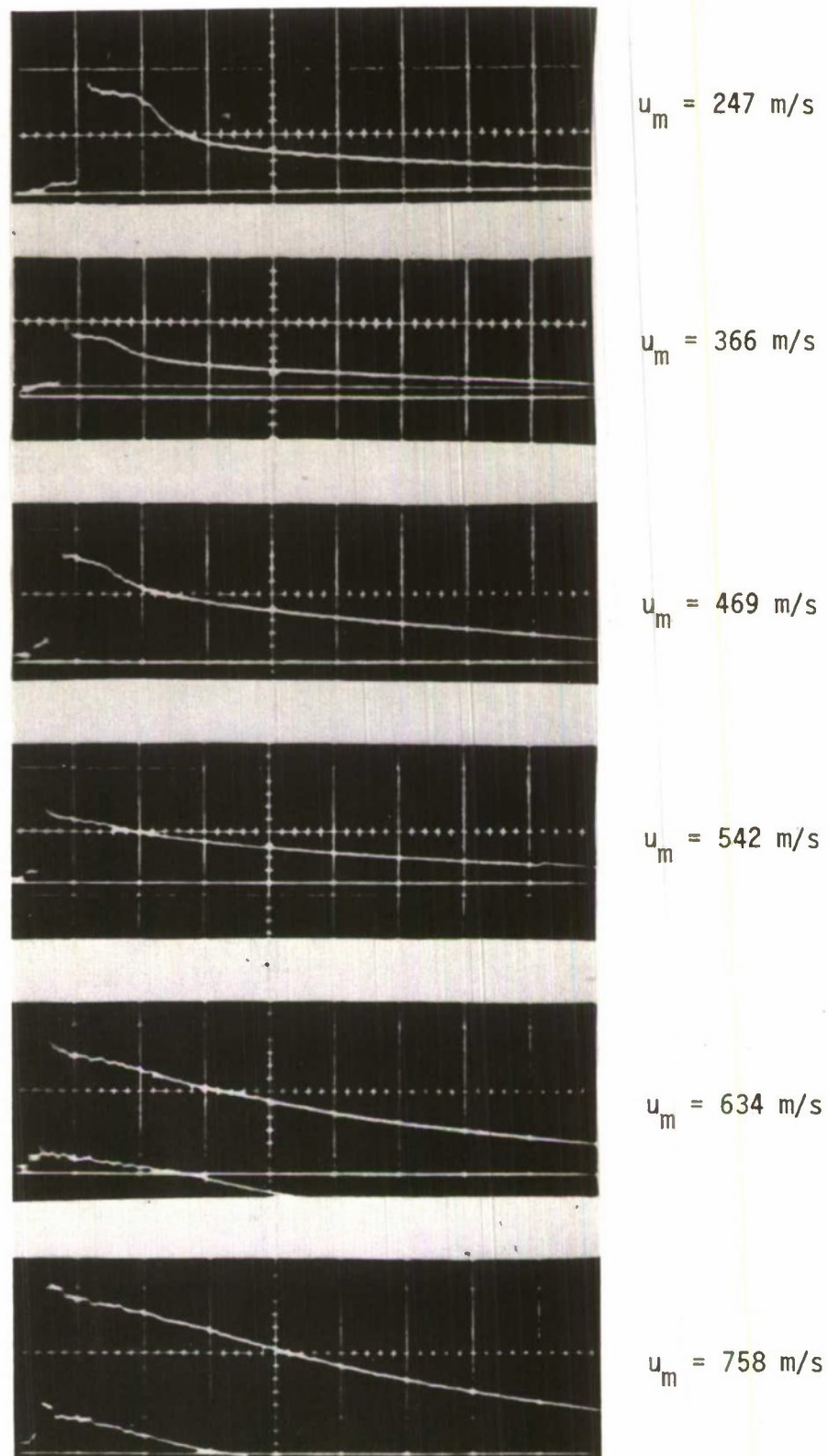


Figure 8. Pressure Traces from Gauge 3 at Various Launch Velocities

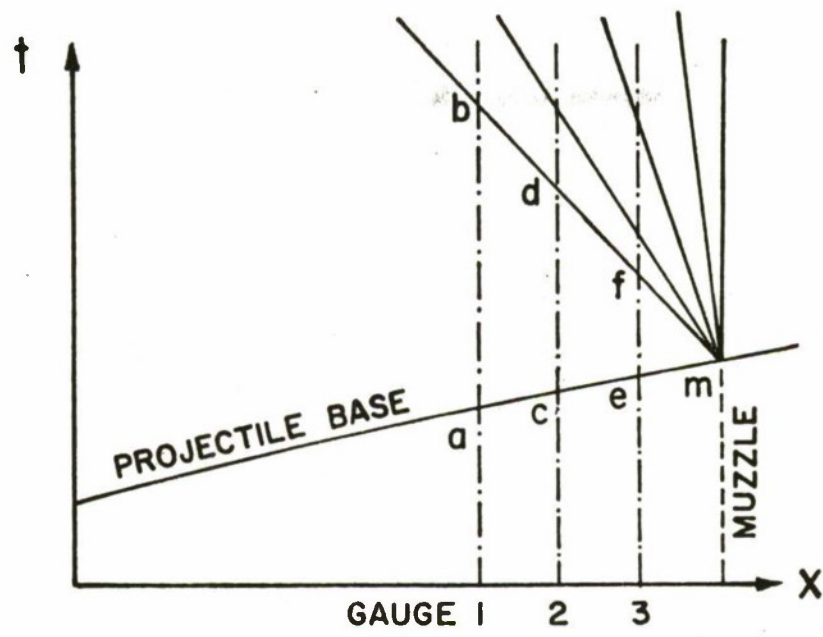


Figure 9. Wave Diagram of Muzzle Flow

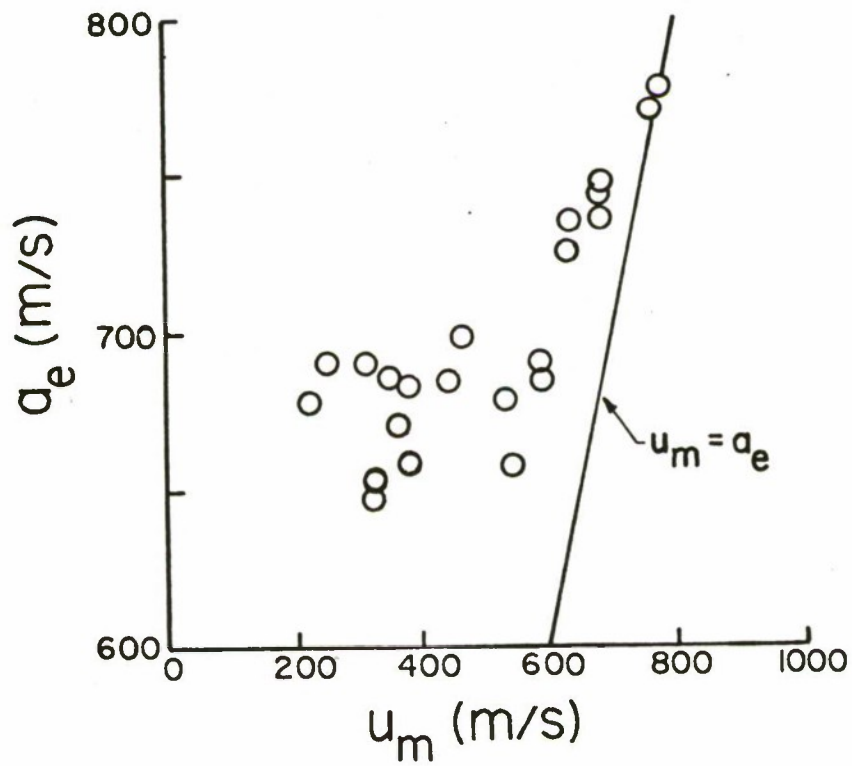


Figure 10. Measured Speed of Sound

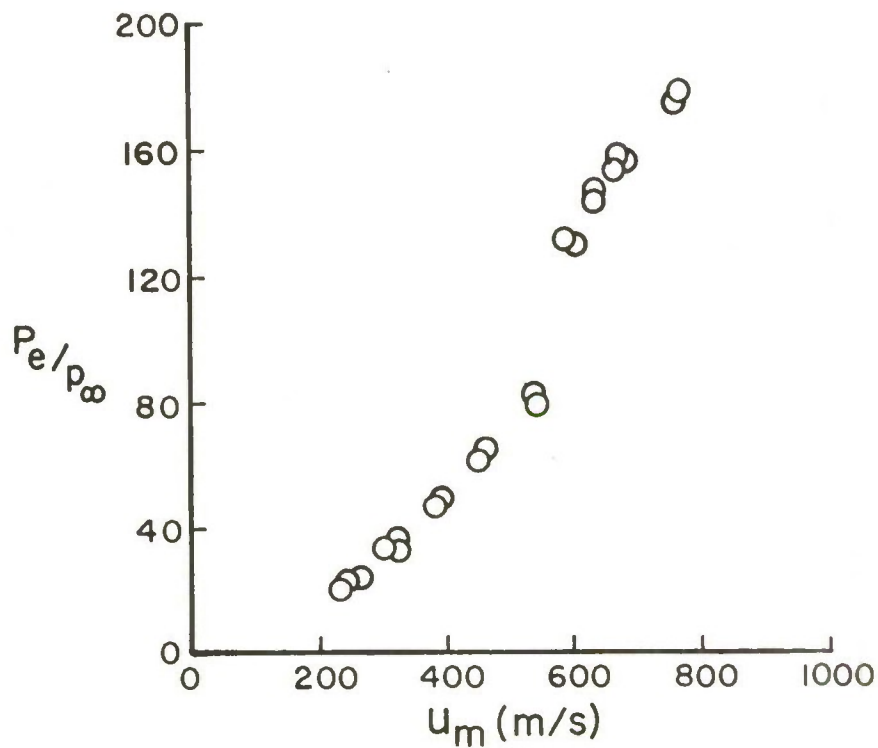


Figure 11. Measured Exit Pressure

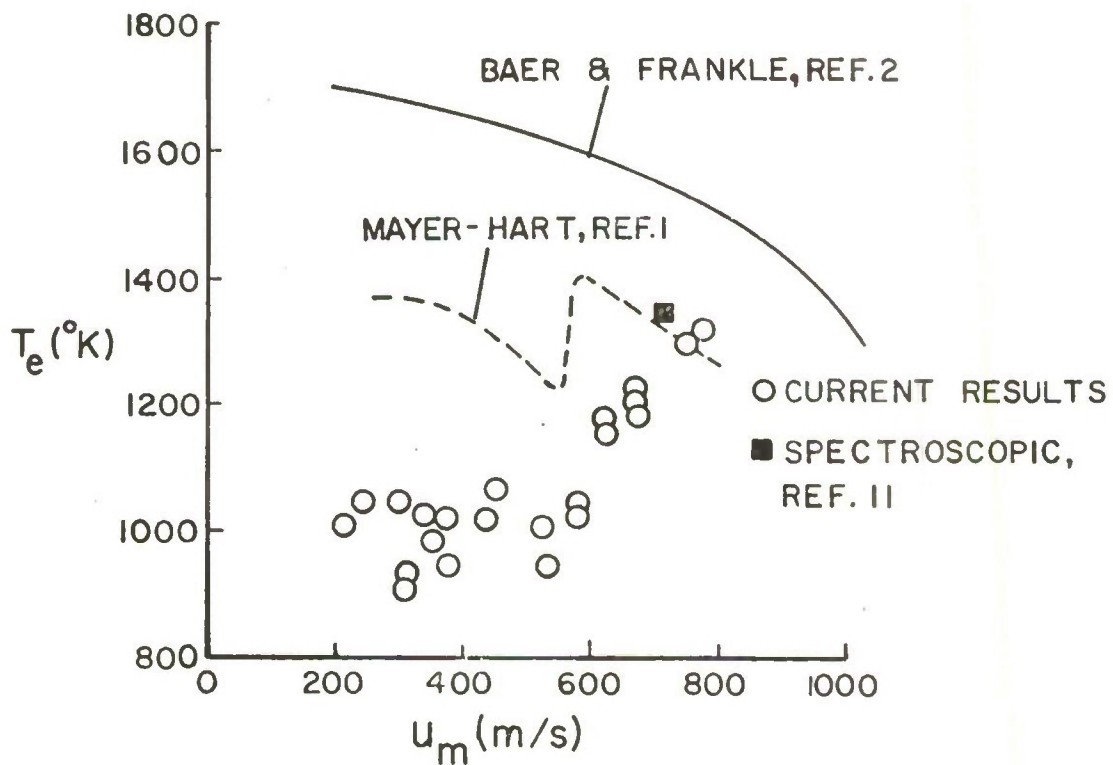


Figure 12. Measured Exit Temperature

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LIST OF SYMBOLS

a	speed of sound
A	area
C_v	specific heat at constant volume
c	wave speed
\dot{E}	rate at which energy passes a given point
p	pressure
R	gas constant
t	time
T	temperature
u	velocity
x	axial coordinate
γ	ratio of specific heats

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